

Conformal bootstrap approach to critical phenomena: A status report

Slava Rychkov



Institut des Hautes Études Scientifiques,
Bures-sur-Yvette

slava.rychkov@protonmail.com



SIM NS
FOUNDATION

STATPHYS 29 - Firenze - 15.07.2025

Two points of view on Conformal Field Theories

1) IR fixed points of microscopic theories
described by a Lagrangian or a lattice model

conformal symmetry
is emergent

THIS TALK:

2) Defined algebro-analytically as
systems of correlators of local operators
satisfying the operator product expansion

conformal symmetry
is built in

= conformal bootstrap

general philosophy	Polyakov 1974
D=2	Belavin, Polyakov, Zamolodchikov 1984 talk by Belavin @ STATPHYS Edinburgh 1983
general D	Rattazzi, S.R., Tonni, Vichi 2008 talk by S.R. @ STATPHYS Seoul 2013

Conformal Bootstrap Algorithm

(2008-present)

A physical system (Lagrangian, lattice model...)
giving rise to a unitary CFT in D or $(D-1)+1$ dimensions

- Symmetry
- A guess about operator spectrum (e.g. relevant scalars)

Split Operators = **Low** + **High**

E.g. **Low**= relevant scalars, symmetry currents,
stress tensor,...

High = Everyone else

Crossing equation:

$$\langle \mathcal{O}_i \mathcal{O}_j \mathcal{O}_k \mathcal{O}_l \rangle = \sum_r \text{ "s-channel" } = \sum_{r'} \text{ "t-channel" }$$

OPE coeffs

“s-channel”

“t-channel”

$i, j, k, l \in \mathbf{Low}$

Exchanged $r, r' = \mathbf{Low+High}$

- Equation for functions of $x_1, x_2, x_3, x_4 \in \mathbb{R}^D$
 $\implies$ reduced by conformal symmetry to $(z, \bar{z}) \in \mathbb{C}$
- Expand to order Λ around “half-way” point $z = \bar{z} = 1/2$ where
 - both sides converge exponentially fast
 - terms have good positivity properties

Rattazzi, **SR**, Tonni, Vichi
 JHEP **0812** (2008) 031

Expansion order $\Lambda \lesssim 1000$ depending on your computer

“Scan” over the **Low** operators and their OPE coefficients
to find **allowed regions**

(defined as regions where *some* **High** operators exist
so that crossing holds, i.e. marginalizing over **High**)

CFT parameter space

“Guess” from
microscopics



Allowed

Smart ways of scanning

(Ning Su)

- cutting surface algorithm
- navigator function, ...

— Analysis is rigorous because of positivity (LP, SDP)

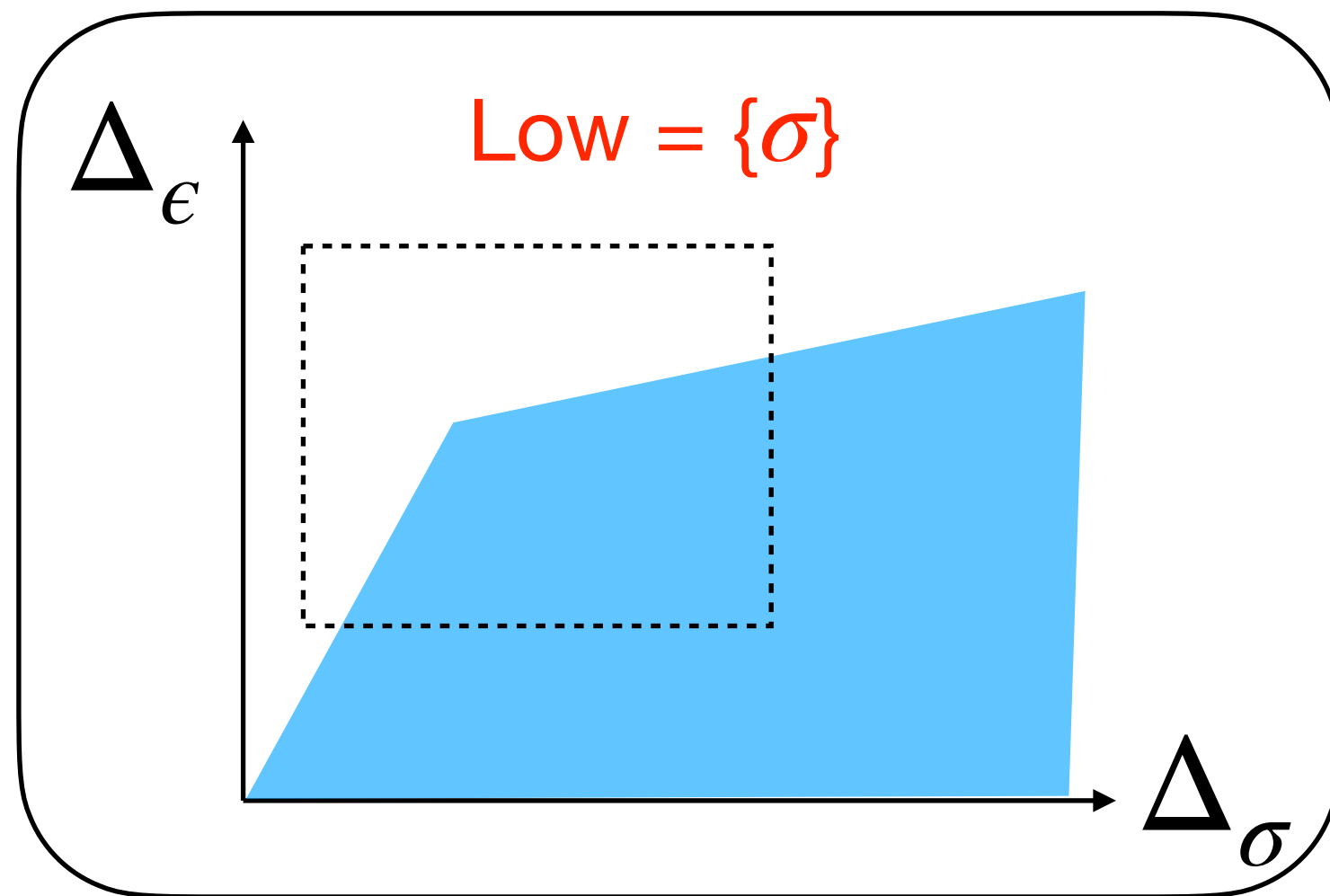
Kos, Poland, Simmons-Duffin 2014
Simmons-Duffin 2015

— Allowed regions always shrink imposing more constraints

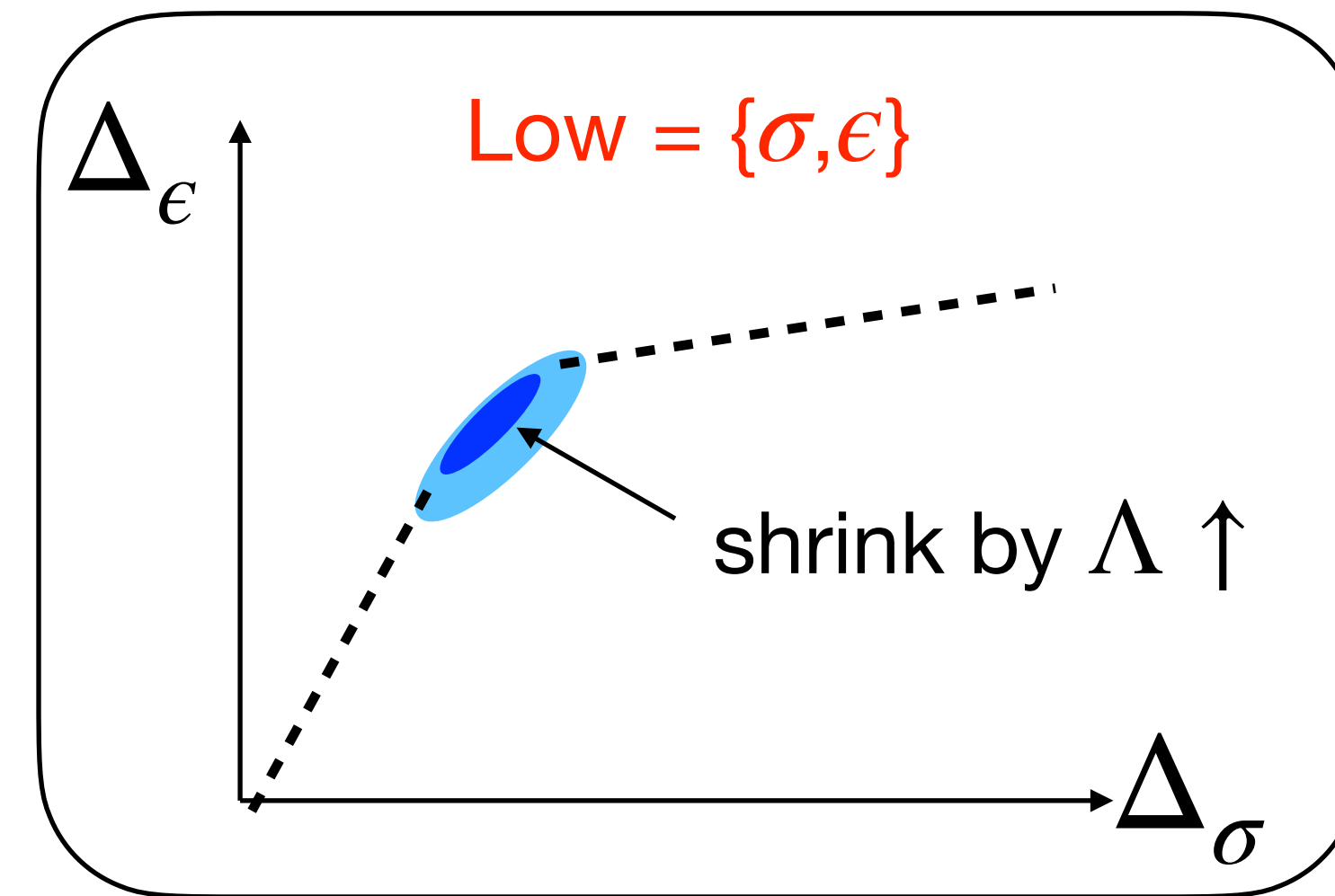
(higher Δ , more **Low** operators)

Highlight 1: 3D Ising CFT

$\mathbb{Z}_2$ symmetry. Two relevant scalars: σ ($\mathbb{Z}_2$ odd), ϵ ($\mathbb{Z}_2$ even)



El-Showk, Paulos, Poland,
SR, Simmons-Duffin, Vichi 2012, 2014



Kos, Poland, Simmons-Duffin 2014
Kos Poland Simmons-Duffin Vichi 2016

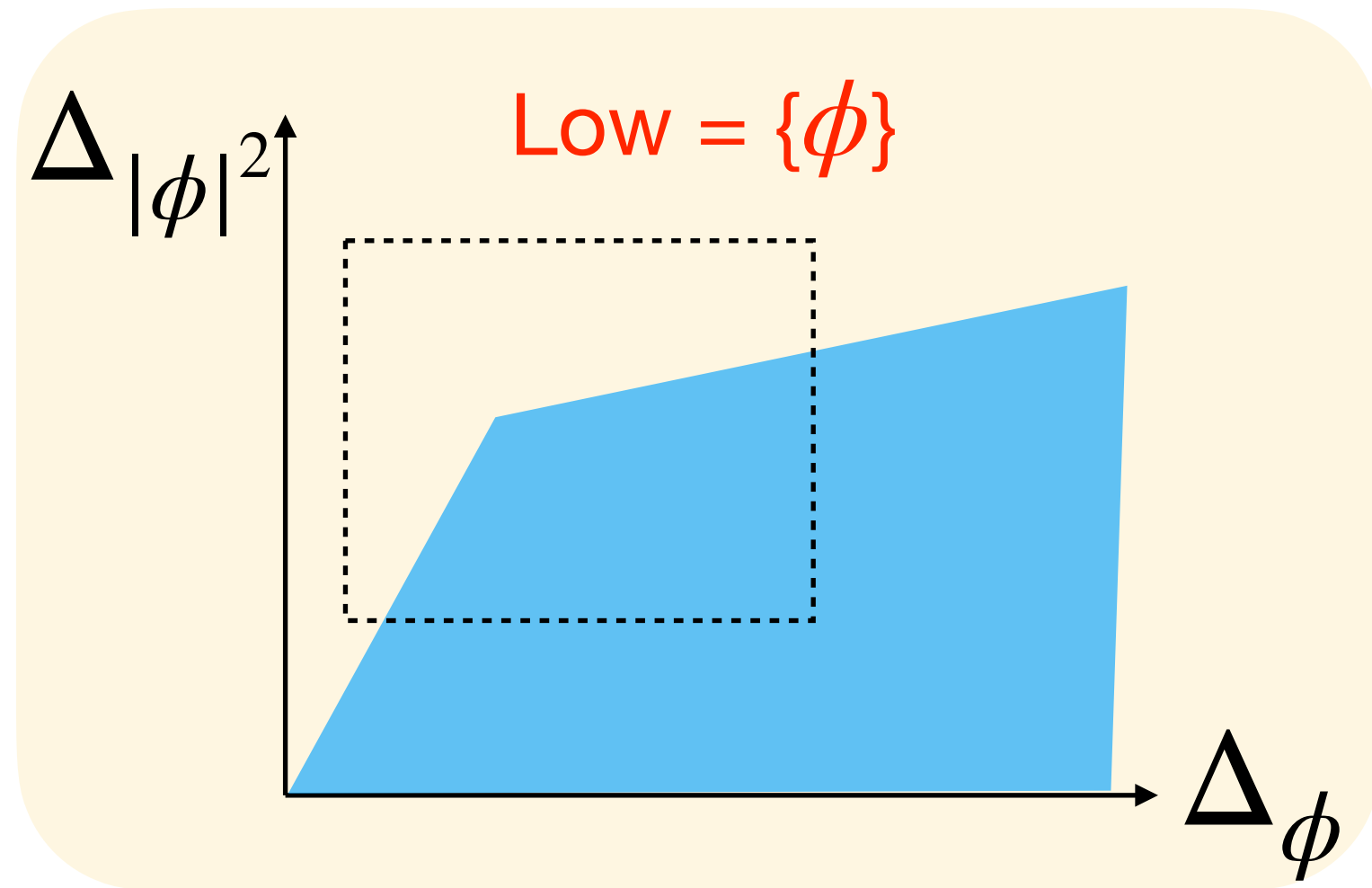
$$\Delta_\sigma = 0.5181489(10)$$
$$\Delta_\epsilon = 1.412625(10)$$
$$\sim 10^{-5} \text{ accuracy}$$

$$\text{Low} = \{\sigma, \epsilon, T_{\mu\nu}\}$$

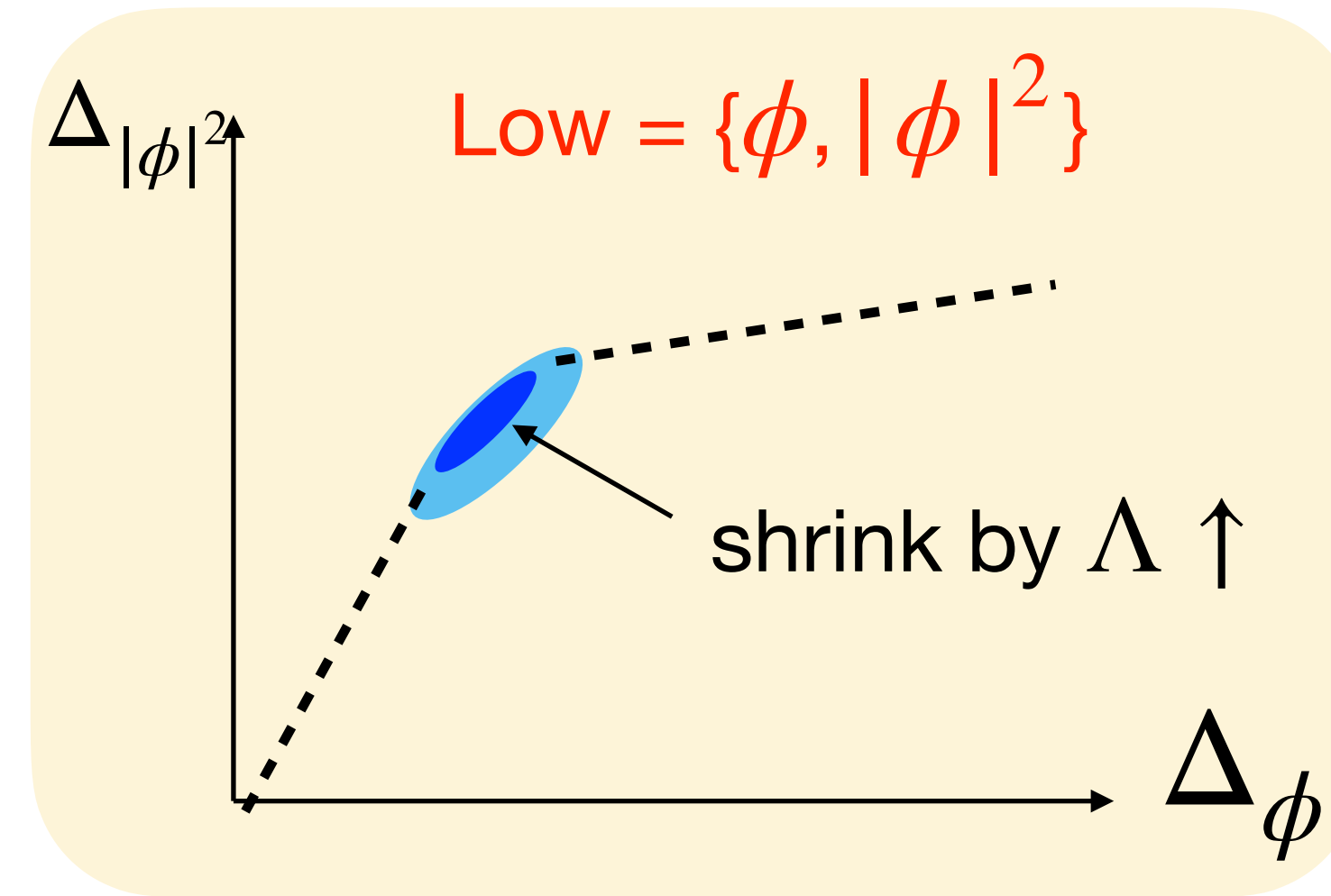
See NEXT TALK Rajeev Erramilli

Highlight 2: 3D XY model CFT

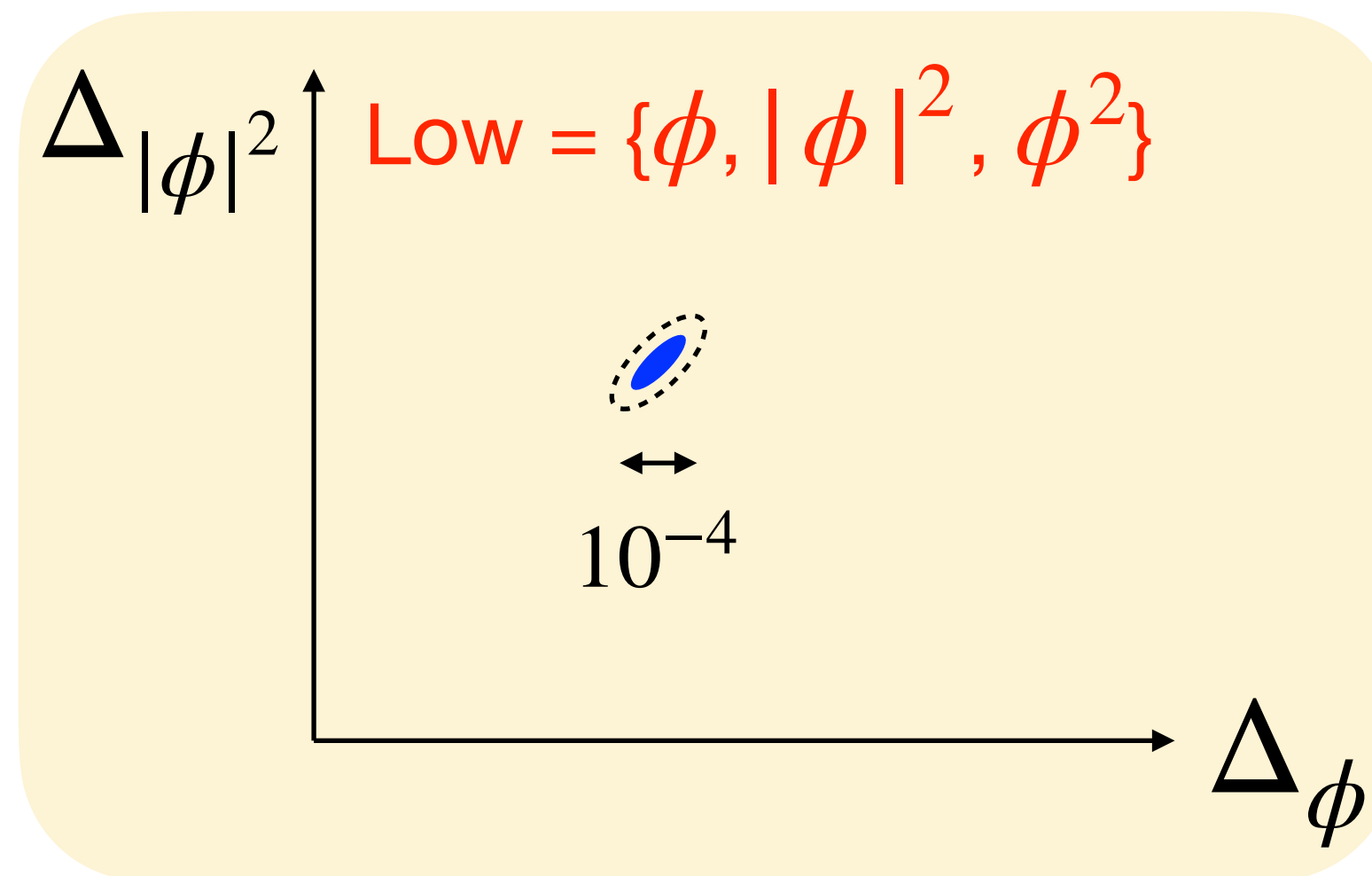
O(2) symmetry. 4 relevant scalars: ϕ , $|\phi|^2$, ϕ^2 , ϕ^3



Kos, Poland, Simmons-Duffin 2013



Kos, Poland, Simmons-Duffin, Vichi 2015



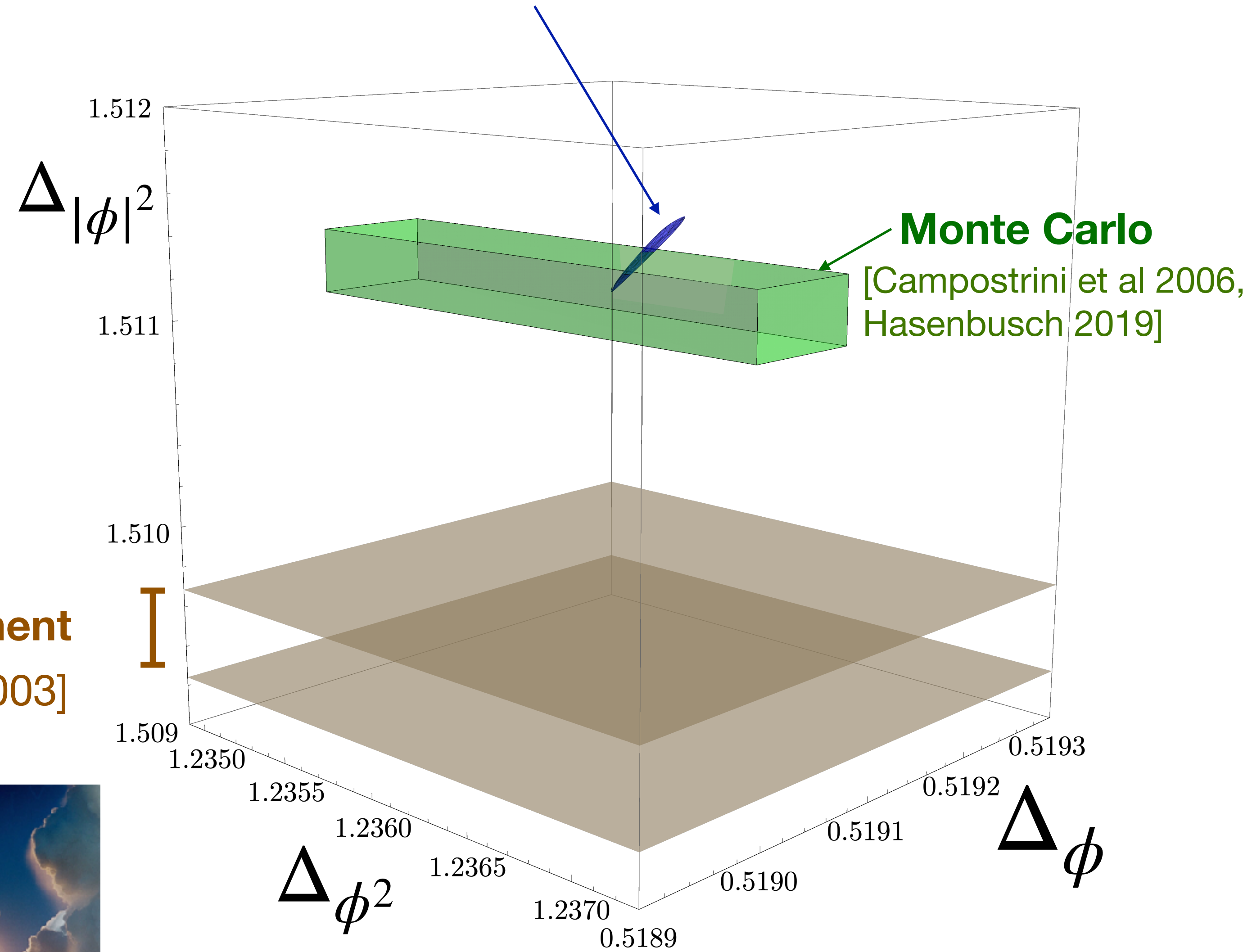
$$\Delta_\phi = 0.519088(22)$$

$$\Delta_{|\phi|^2} = 1.51136(22)$$

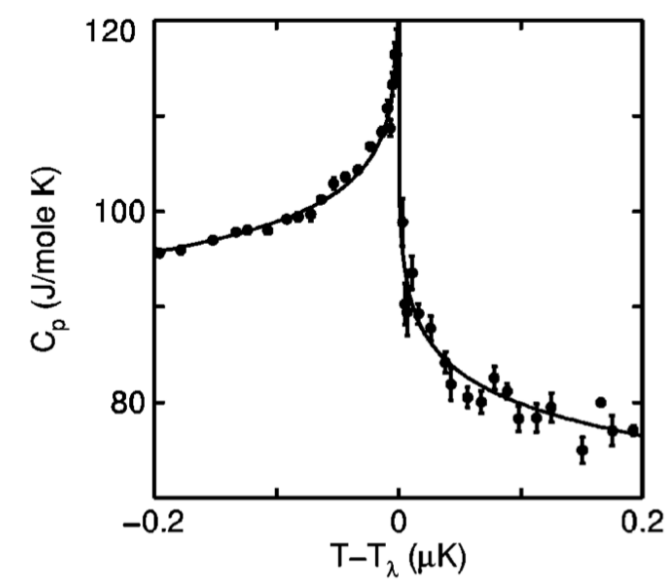
Chester, Landry, Liu, Poland,
Simmons-Duffin, Su, Vichi 2020

Bootstrap

[Chester, Landry, Liu, Poland, Simmons-Duffin, Su, Vichi, 2019]

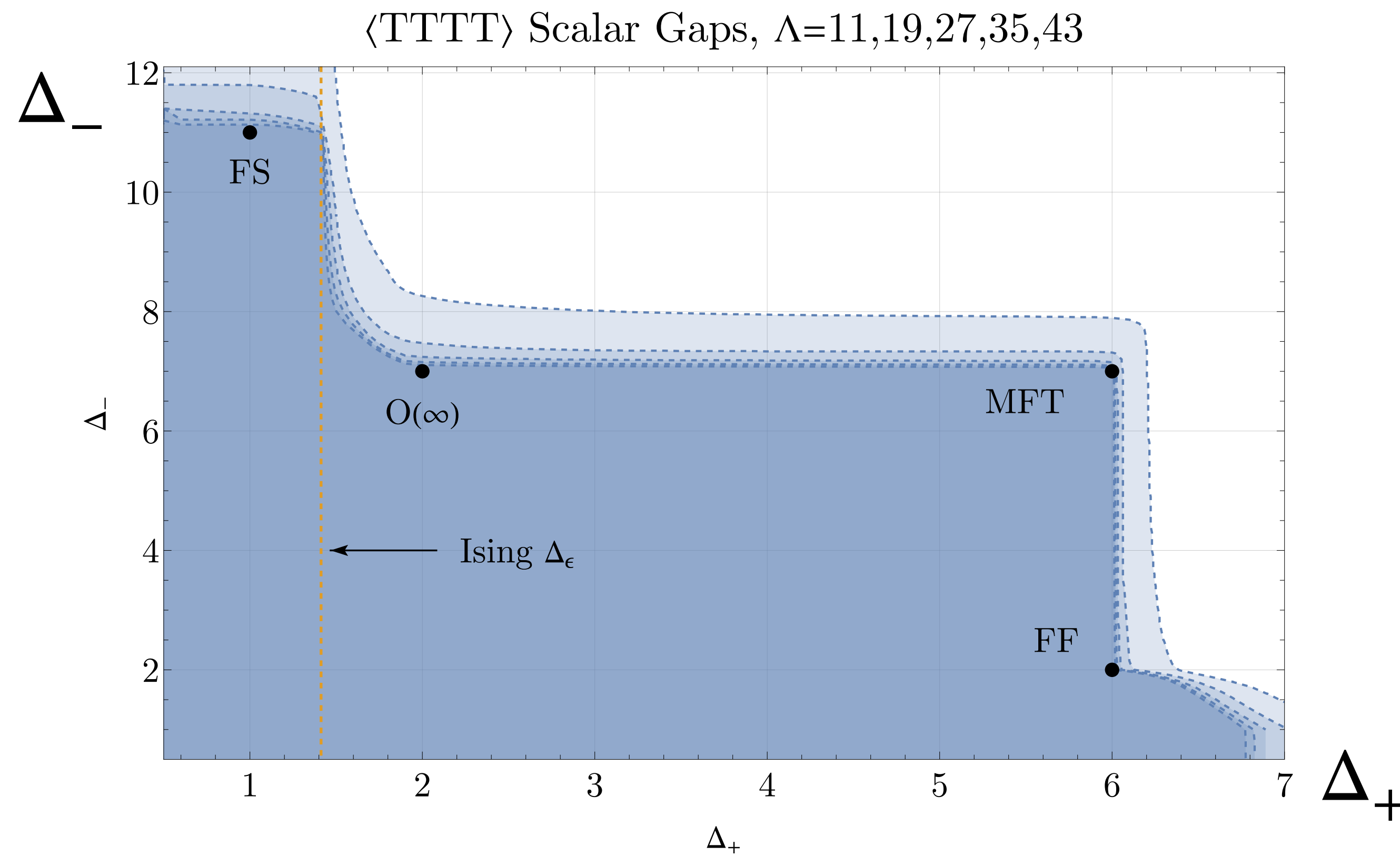


He-4 experiment
[Lipa et al, 2003]



Highlight 3: Space of all theories

- Take ANY **P-invariant 3D CFT**, P =spatial parity
- **Low**= $\{T_{\mu\nu}\}$ stress tensor
- Bounds on $\Delta_{\pm}$ - dimensions of lowest parity-even/odd scalars in TxB OPE



plot courtesy of
Rajeev Erramilli

Dymarsky, Kos, Kravchuk, Poland, Simmons-Duffin'2017

Selected open problems



1) Uniqueness/Non-existence problems

- show that CFTs describing most famous universality classes are unique
- show that there is no CFT when phase transition is 1st order

2) Bootstrapping gauge theories

- isolate into bootstrap islands IR fixed points of gauge theories

3) 'Large Δ problem' $\Rightarrow$ "analytic functional bootstrap"?

- speed up convergence of bootstrap computations using a smarter basis of functionals

Open Problems - BACK UP

See also the recording:

<https://video.desy.de/video/wpc-theoretical-physics-symposium-2025-slavarychkov/5b868b5815ca6e130168f5c2e556aa4d>

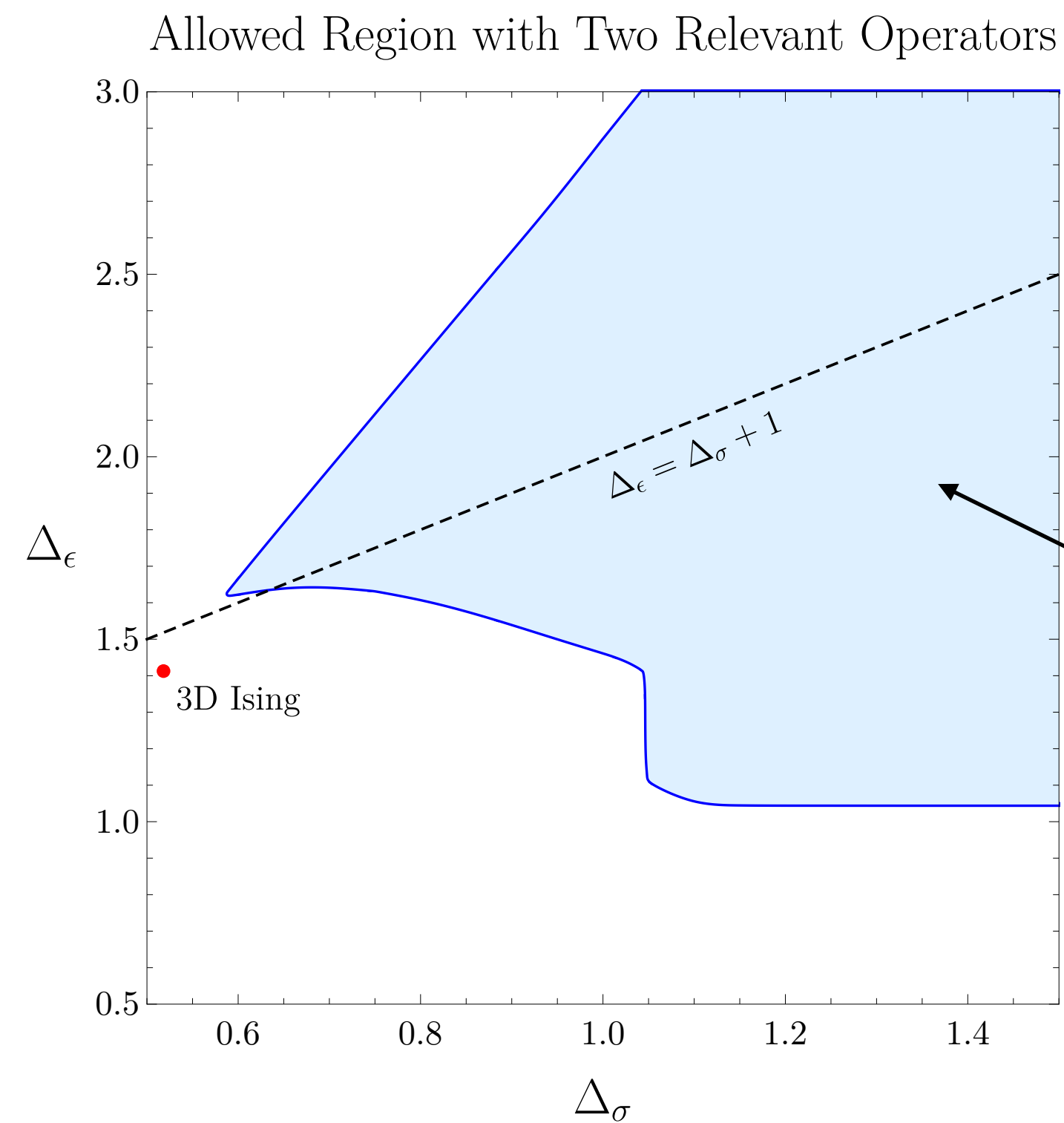
1A) Uniqueness problem

Is 3D Ising CFT unique?

Experiments suggests **yes**

(if not we'd see Ising magnets/liquid-vapor critical points with other exponents)

Can we show this rigorously via bootstrap?



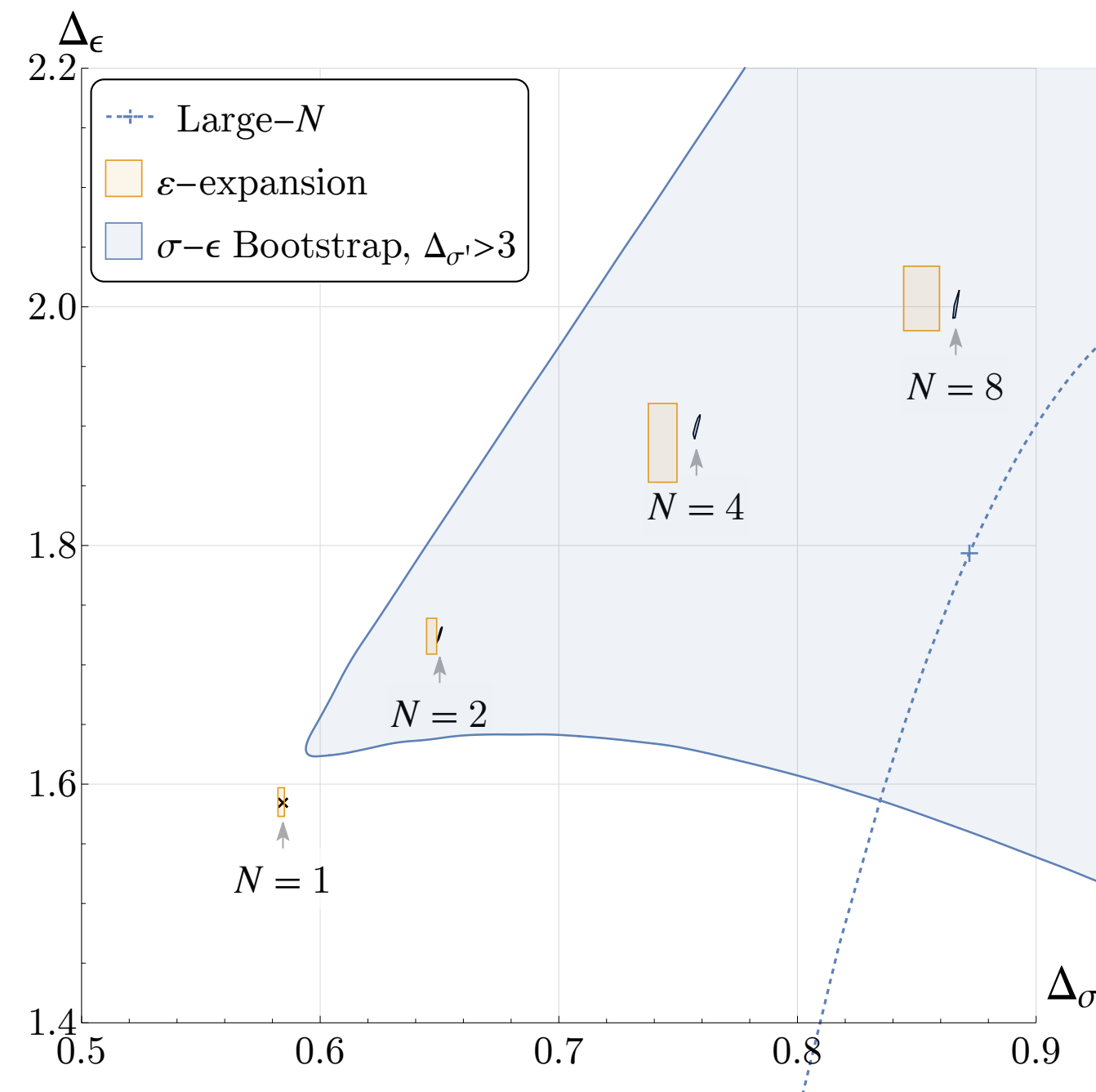
Atanasov, Hillman, Poland 2022

What's the meaning
of this region?

Gross-Neveu-Yukawa model masquerading as Ising

$$\mathcal{L}_{\text{GNY}} = -\frac{1}{2}(\partial\phi)^2 - i\frac{1}{2}\psi_i\partial\psi_i - \frac{1}{2}m^2\phi^2 - \frac{\lambda}{4}\phi^4 - i\frac{g}{2}\phi\psi_i\psi_i$$

The $O(N)$ Gross-Neveu-Yukawa Archipelago



Under spatial parity

$$P : \psi\psi \rightarrow -\psi\psi, \quad \phi \rightarrow -\phi$$

Erramilli, Iliesiu, Kravchuk, Liu, Poland, Simmons-Duffin, Su, Vichi 2023

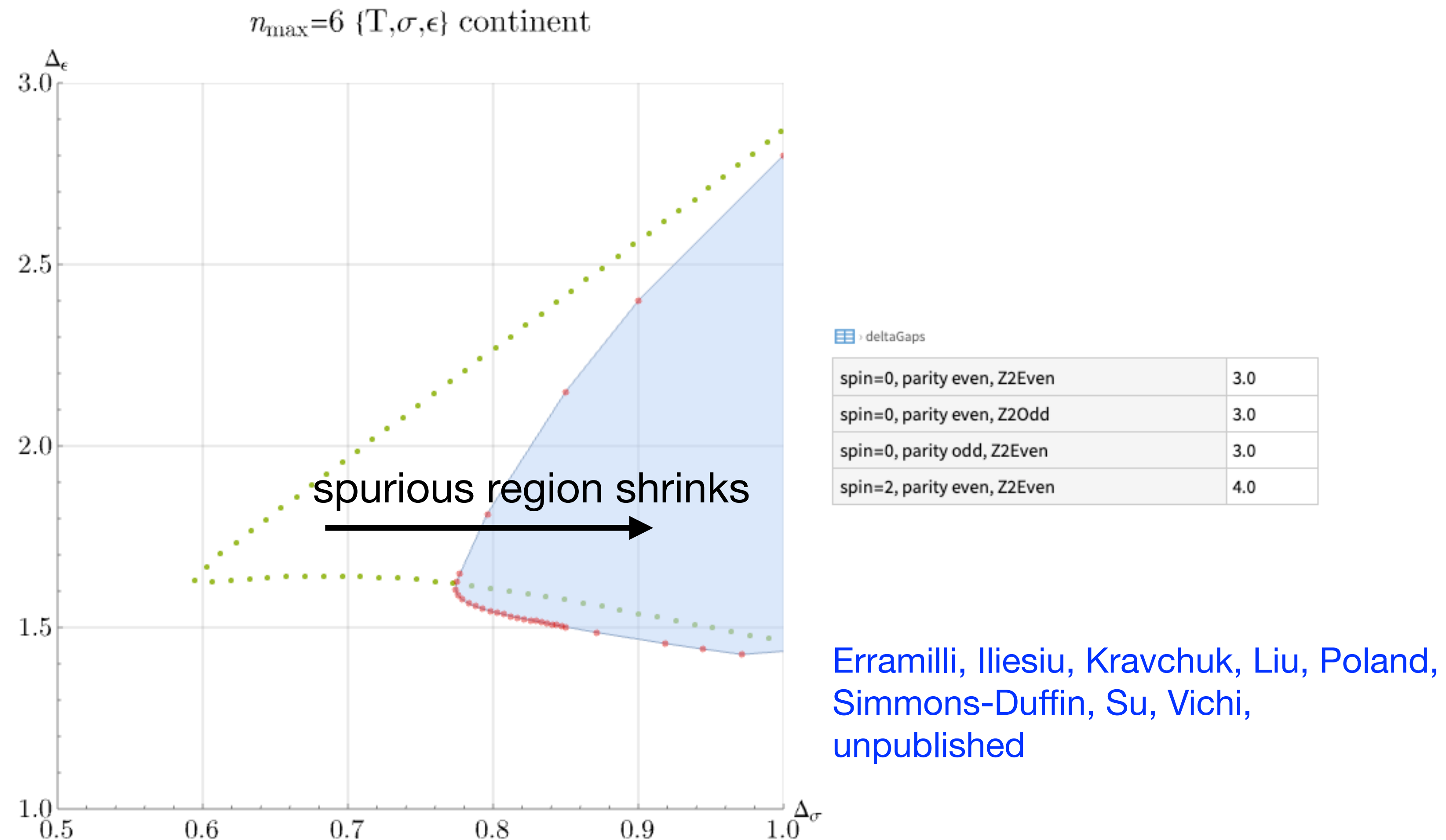
- For $\text{Low} = \{\phi, \phi^2\}$, P is indistinguishable from Ising $\mathbb{Z}_2$

- May be distinguished for $\text{Low} = \{\phi, \phi^2, T_{\mu\nu}\}$

$$T \times T \supset \phi \quad \text{in GNY but not in Ising}$$

Ising and GNY may be distinguished for $\text{Low} = \{\phi, \phi^2, T_{\mu\nu}\}$

$T \times T \supset \phi$ in GNY but not in Ising



Open problem: Can the spurious region be eliminated altogether?

1B) Non-existence problem

For some models experiments and Monte Carlo suggest 1st order transition

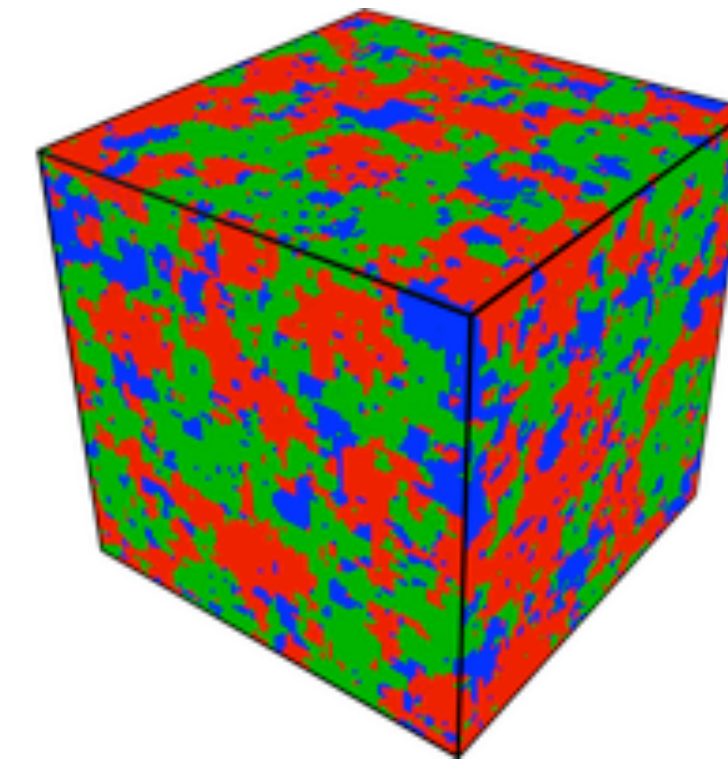
But one can never quite exclude 2nd order in a slightly modified model.

A **proof** can be obtained by showing that there is no CFT with requisite symmetry.

Simplest case: 3-state Potts model in $D=3$

Lattice Monte Carlo: correlation length $\xi \sim 10$

[Janke, Villanova 1997]



Bootstrap open problem:

Show that there is no unitary 3D CFT

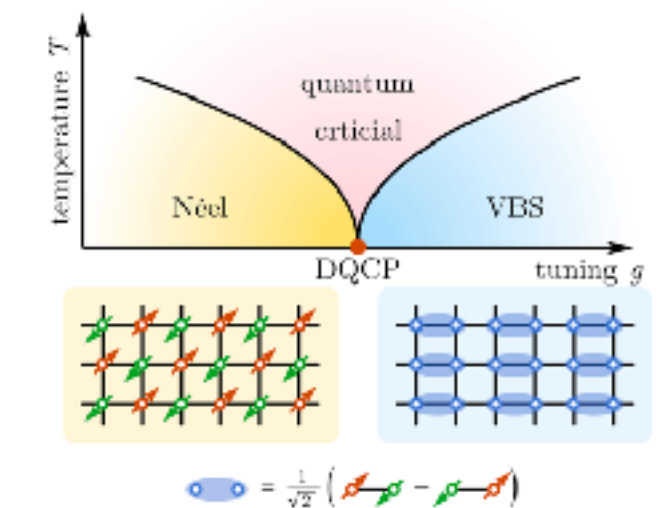
- with S_3 global symmetry
- a single relevant singlet scalar
- one (or more) scalars in the fundamental of S_3

2) Bootstrapping 3D conformal gauge theories

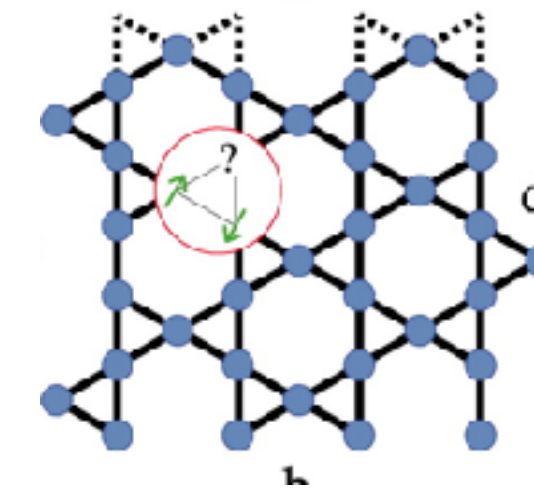
QED3 (bosonic/fermionic) = 3D U(1) Maxwell field + N_f bosons/fermions

Global symmetry: $G \simeq SU(N_f) \times U(1)_{top}$

Bosonic QED3 $N_f = 2$ $\longleftrightarrow$ Deconfined Quantum Critical Point



Fermionic QED3 $N_f = 4$ $\longleftrightarrow$ Dirac Spin Liquid



Herbertsmithite



Can we bootstrap these CFTs and determine N_f^* ?

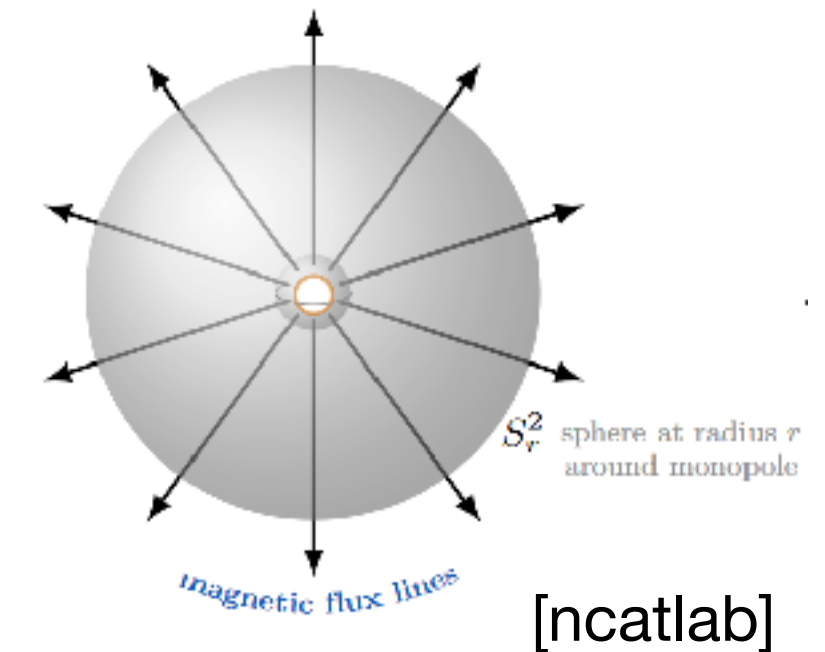
Features of QED3

- Physical CFT operators are gauge invariant combinations of elementary fields

$$\bar{\psi}_i \psi_j, \quad \phi_i^* \phi_j, \dots \quad \text{---} \psi, \phi$$

- There are also “monopole operators” charged under $U(1)_{top}$

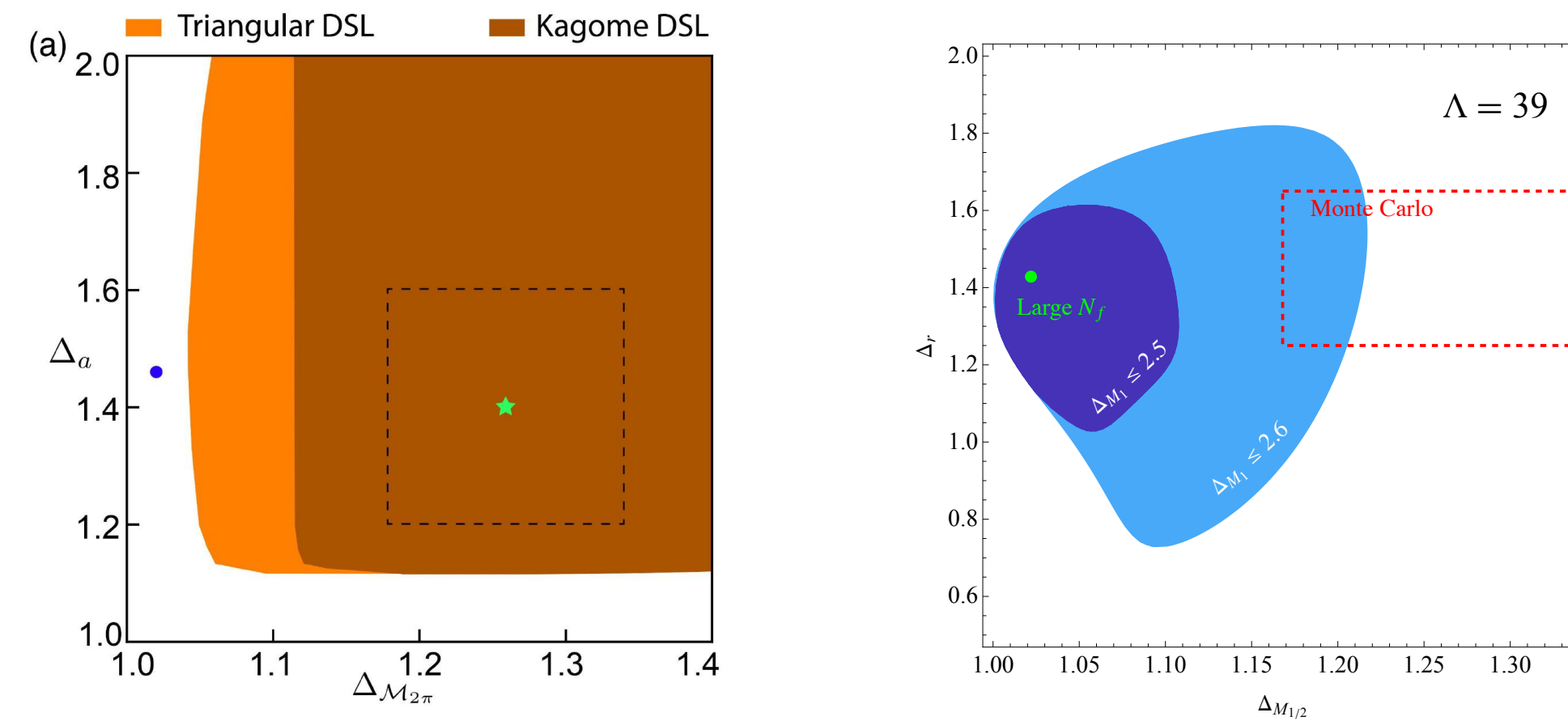
$$\Delta_q \propto N_f$$



Difficulties:

- These operators are heavier than the lightest scalars in scalar/fermionic CFTs
 - expect slower convergence as $\Lambda \uparrow$ (cf “large Δ problem”)
- How to distinguish from “QCD3” theories where the gauge group is $U(N_c)$?

Various bootstrap bounds on QED3 were derived but these CFTs were not yet isolated into small cl



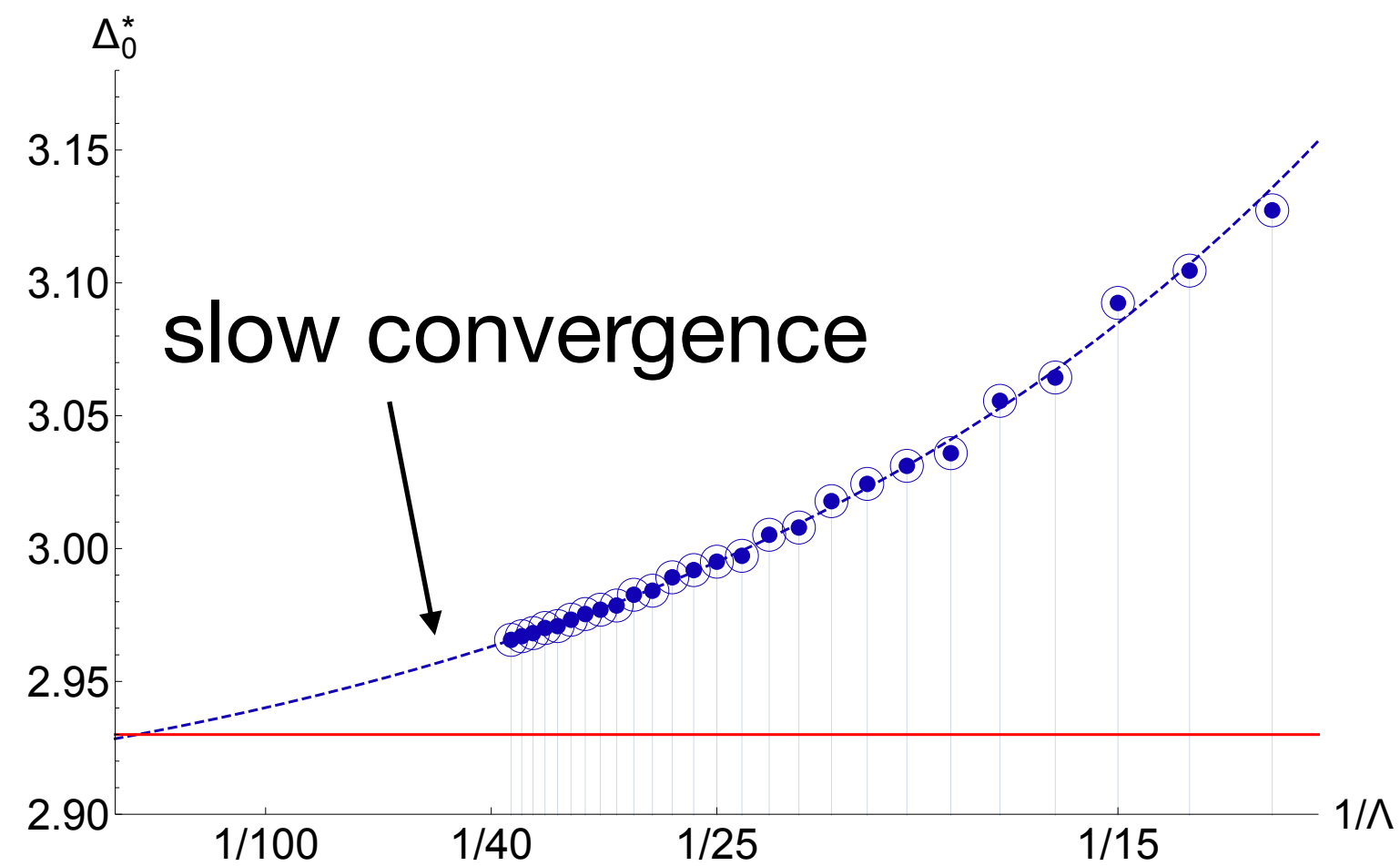
see [SR, Su RMP 2024]
for a discussion

[Reehorst, Refinetti, Vichi 2020, He, Rong, Su 2021]
identified gaps in the operator spectrum (“decoupling operators”)
which could help distinguish QED3 from QCD3

(color indices allow for more antisymmetrization)

3) Large Δ problem

= Slow $\Lambda \uparrow$ convergence for bounds on correlators of large- Δ operators



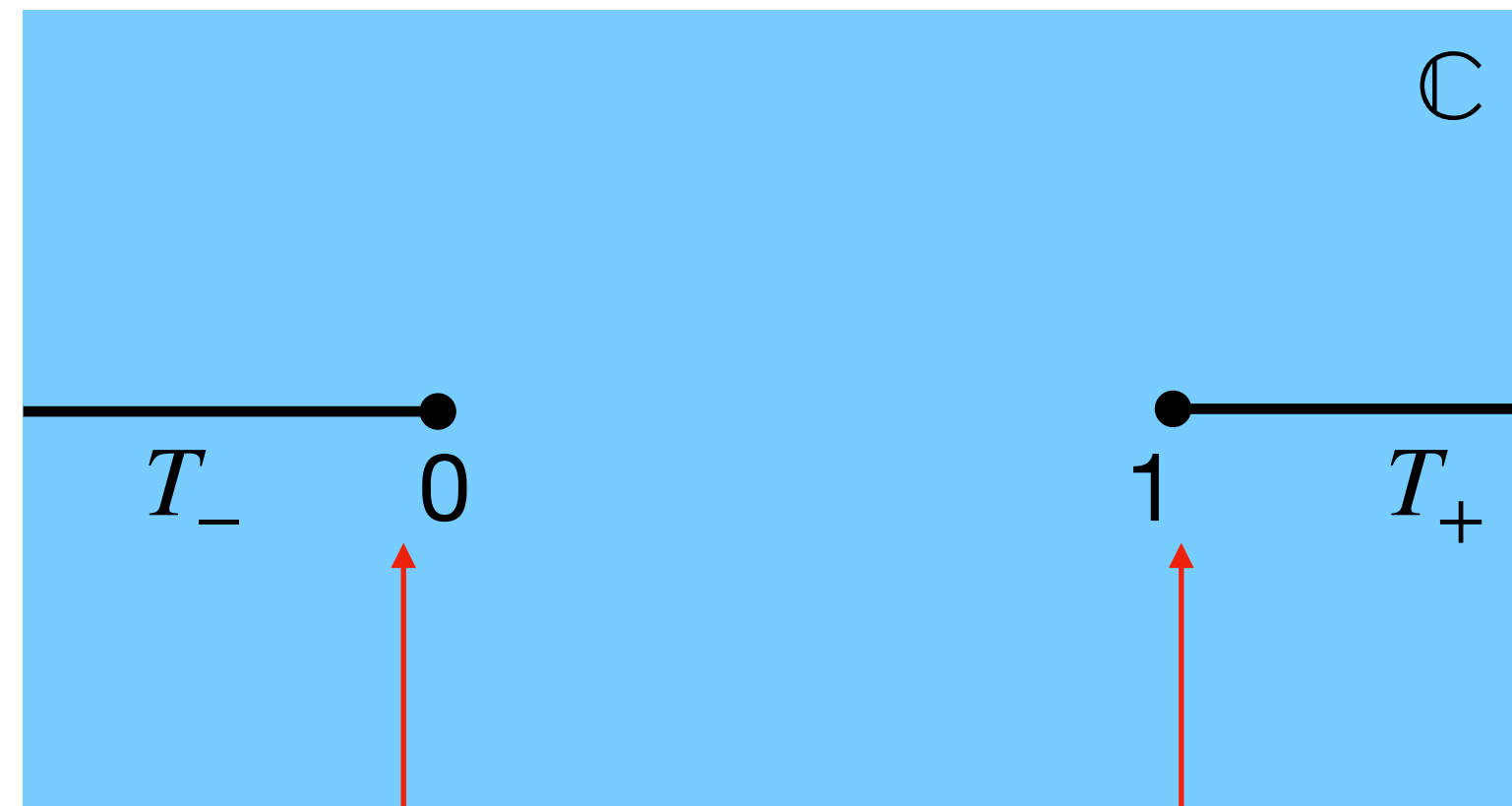
Upper bound on 1st unprotected scalar
for $\mathcal{N} = 4$ SCFT with $c=3/4$
(Konishi in SYM with $SU(2)$ gauge group)

Beem, Rastelli, van Rees 2016

From 4pt function of protected scalar in **20'** of “large” dimension $\Delta = 2$

Origin of large Δ problem

CFT 4pt function $\langle \mathcal{O}_\Delta \mathcal{O}_\Delta \mathcal{O}_\Delta \mathcal{O}_\Delta \rangle$ is analytic for $z, \bar{z} \in \mathbb{C} \setminus (T_+ \cup T_-)$



$$1/(z\bar{z})^\Delta$$

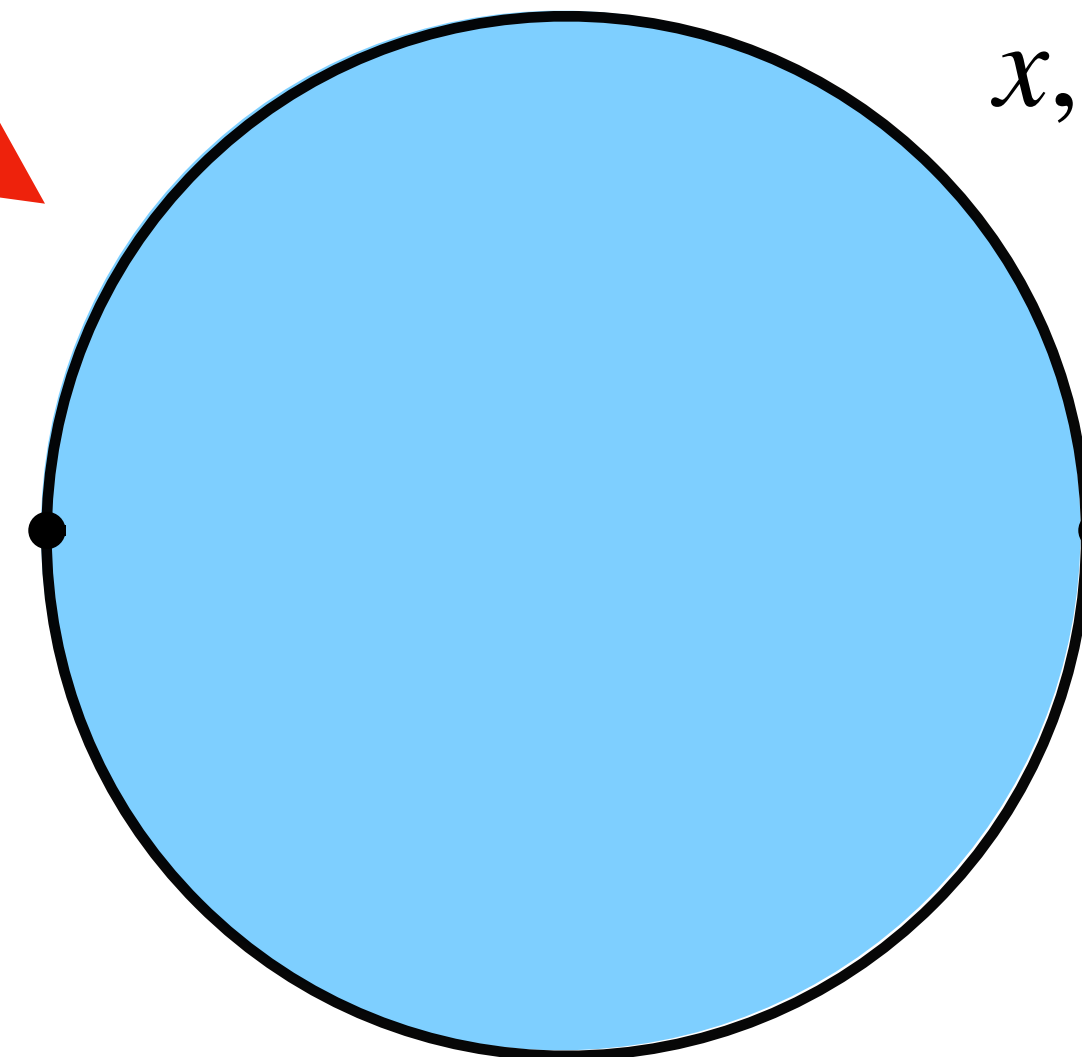
$$1/[(1-z)(1-\bar{z})]^\Delta$$

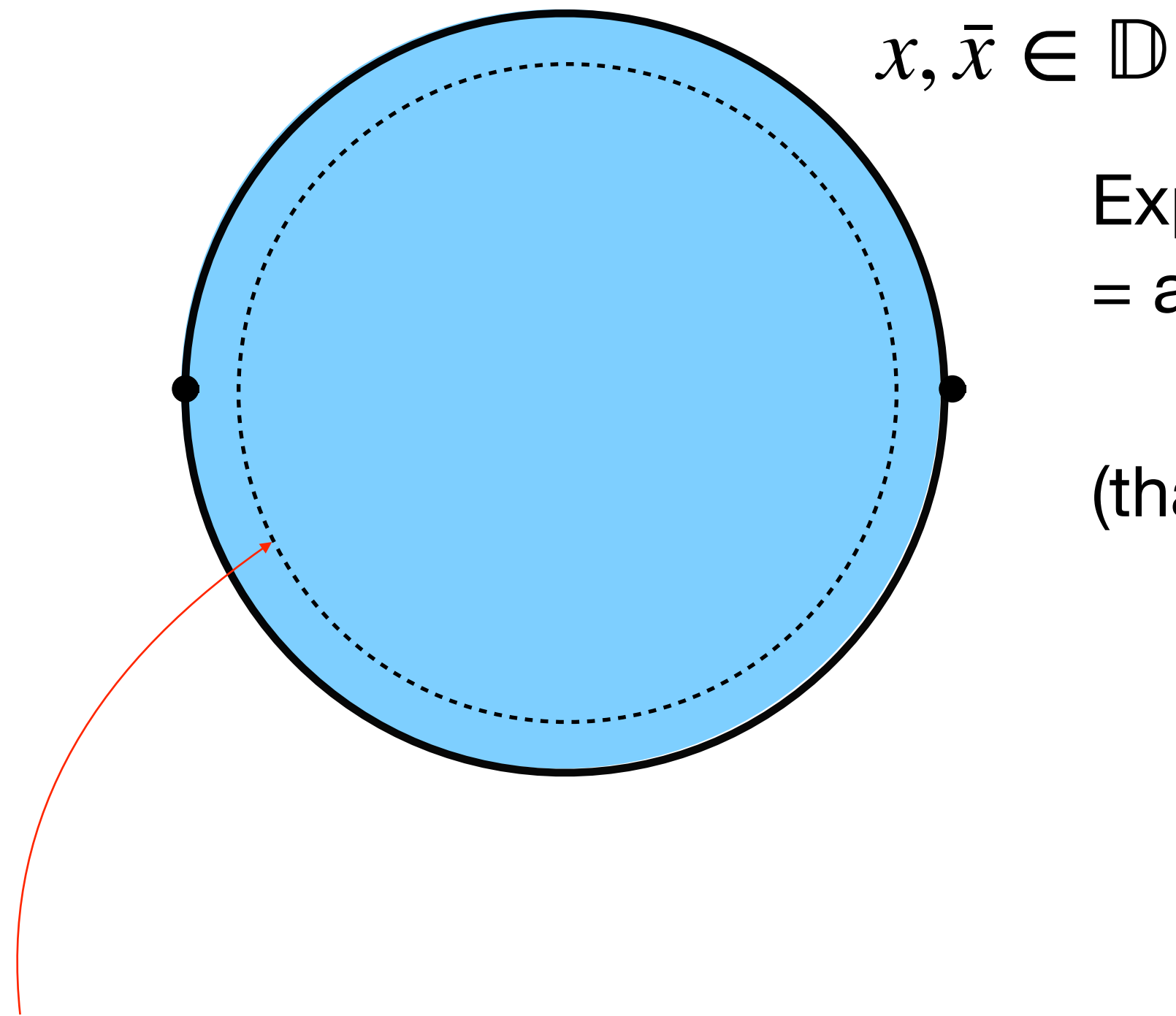
s-channel analytic in $\mathbb{C} \setminus T_+$

t-channel analytic in $\mathbb{C} \setminus T_-$

conformal map

$$x, \bar{x} \in \mathbb{D}$$





Expand crossing equation around $x = \bar{x} = 0$
 = act on it with “derivative functionals”

$$\partial_x^n \partial_{\bar{x}}^m \big|_{x=\bar{x}=0} \quad n + m \leq \Lambda$$

(that’s standard way since our 2008 work)

The largest class of functionals given the analyticity domain
 can be obtained as contour integrals pushed to the boundary.

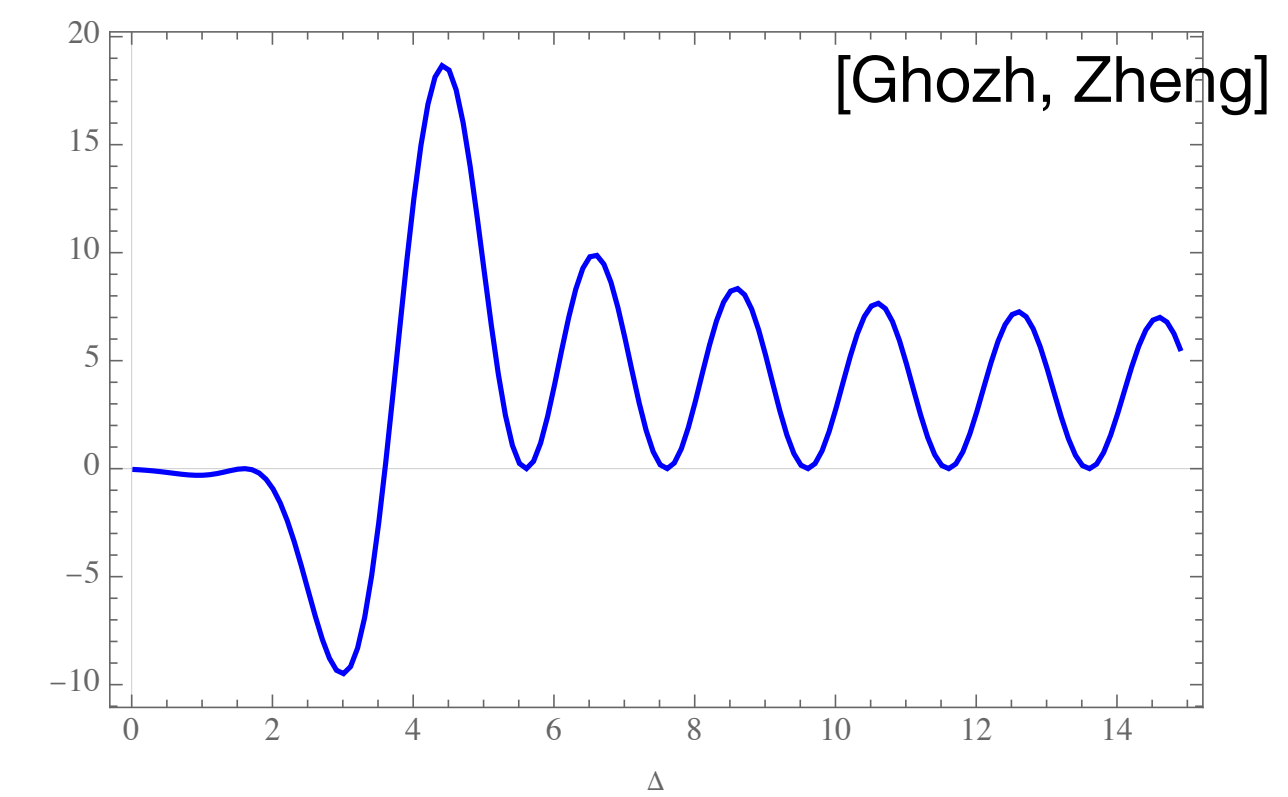
“analytic functionals” Mazac 2016

Mazac, Paulos 2018

Such functionals can be expanded in “derivative functionals”
 but convergence becomes slow for large Δ because of s,t-channel sing’s

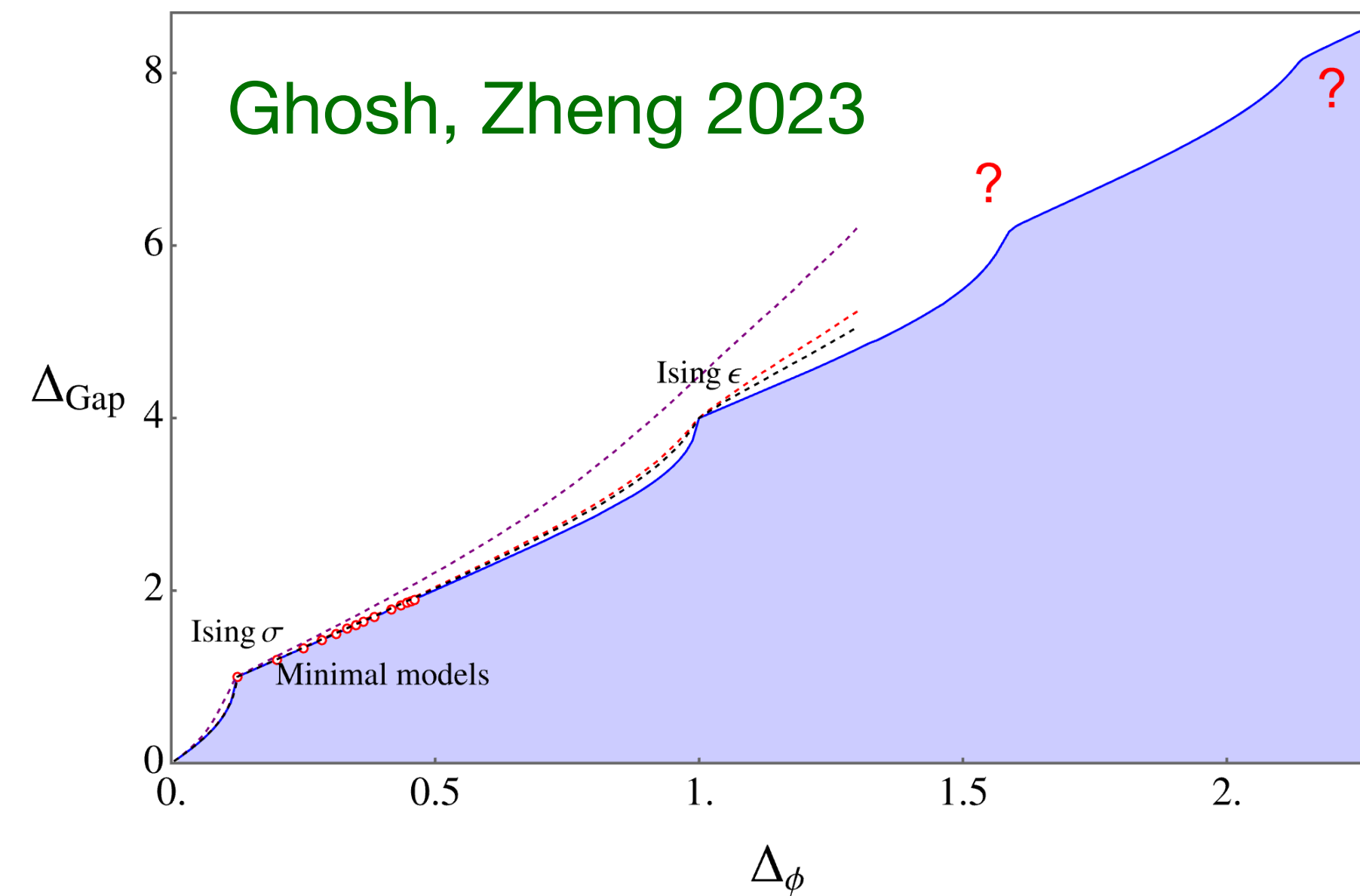
Analytic functionals

Have various magic properties,
give exact solutions to some max gap problems



More generally, *a faster-convergent basis for numerical bootstrap calculations*
Paulos, Zan 2019; Ghosh, Zheng 2023

D=2:



D=3 - works (Ghosh, Zheng 2023) but need more efficient implementation

A future of numerical bootstrap?